

Seeing is Sensing - External Observation of *In Vivo* Biological Gradient Field by Tracking Nanoswimmers

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Abstract—This work aims to demonstrate through computational analysis that, by monitoring the trajectory of externally manipulable nanoswimmers (NS), the *in vivo* biological gradient field (BGF) interacting with the NS can be indirectly observed. This observability is fundamental to the recently proposed framework of computational nanobiosensing (CONA) for “smart” cancer detection. We first present a novel NS propagation model to emulate the complex and chaotic NS kinetics inside the capillary network. Next, we propose an efficient control method that is able to employ the NS as *in vivo* sensors for the measurement of a specific BGF such as blood viscosity. The proposed method, based on the Linear Quadratic Regulator (LQR), effectively stabilizes the signal-to-noise ratio (SNR) induced by the Brownian motion of NS at a level above 10 dB to enhance the accuracy of viscosity estimation.

Index Terms—*In vivo* sensor, Nanoswimmers, Blood viscosity estimation, Linear Quadratic Regulator.

I. INTRODUCTION

Nanoparticle-mediated drug delivery shows great promise in cancer diagnosis [1]. However, without any external guidance, systemic circulation can only bring 0.7% of the injected nanoparticles to the tumor targets [2]. This low efficiency is a key hurdle for translating nanomedicines from the lab into the clinic. From a computational perspective, this sensing process is simply a “brute-force search”, because the nanoparticles (computing agents) can only “detect” a tumor by enumerating all possible pathways in the complex vascular network [3]. By replacing nanoparticles with externally manipulable nanoswimmers (NS), we have proposed a novel computational nanobiosensing (CONA) framework for “smart” searching of tumors [4]. The presence of a tumor induces changes in some biological gradient fields (BGFs) such as pathological vascular structure, blood velocity and viscosity, pH, enzymatic activity, and homeostatic regulation [3], which will subsequently alter the externally measurable properties of NS such as trajectories and aggregation patterns. Computationally, the BGF could be utilized as the “objective function”; the targeted tumor site corresponds to the “global

optimum”; and the NS are “computing agents”. Thus, NS could cooperate to find the targeted site under the guidance of an external steering field programmed according to specific *in vivo* function optimization algorithms. However, The CONA framework is still hypothetical because BGFs were assumed to be perfectly inferred by measuring the characteristics of NS. The current work aims at exploring the observability of BGF, which is a fundamental problem to be investigated in CONA.

Recent studies have suggested that blood viscosity is a useful BGF parameter for early diagnosis of tumors and other diseases [5]. Furthermore, with the continuous development of nanotechnology, magnetic nanoparticles have been extensively used as NS because of their great promise in both controllability and flexibility. The motion of NS in the stagnant fluidic environment has already been reported in some works. The influence on the swimming speed of bacterial NS in high-viscosity fluids was modelled in [6]. The swimming characteristics of ferromagnetic microsphere chains in fluidic environments with respect to viscosity were described in [7]. However, there are rarely NS-based sensing methods related to the hemodynamic changes, especially in terms of blood viscosity. Also, Some nanorobotic control strategies have been verified in the *in vivo* experiment of mice. All of them are only focused on controllability but are short of researching the BGF observability.

In this work, we propose a novel model to emulate the complex and chaotic kinematics of magnetic-nanoparticles-assembled NS in human vessels. Through analysing the average trajectory of several NS released one by one in the time domain, it is able to solve the inverse problem about hemodynamics so as to obtain the estimation of the relative blood viscosity. To further improve the sensing performance, we present a controller that uses the Linear Quadratic Regulator (LQR) for reducing the noise in the viscosity measurement.

The remaining part of this paper is organized as follows: In section II, the system models and the motion-based viscosity sensing strategy are presented. Section III introduces the concept of location-dependent signal-to-noise ratio (SNR) and propose the LQR control method to improve the sensing performance. Numerical results are presented in Section IV to verify the effectiveness of the proposed method. Finally, some concluding remarks are drawn in section V.

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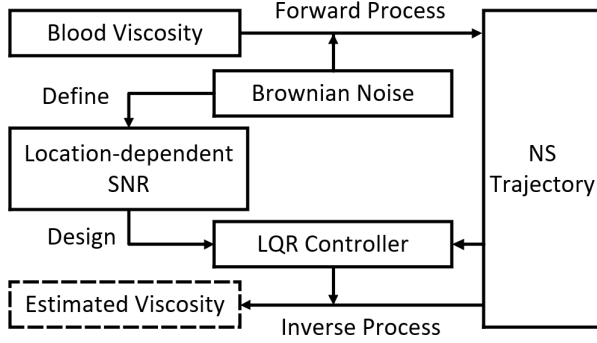


Fig. 1: Framework of the viscosity sensing by tracking the NS.

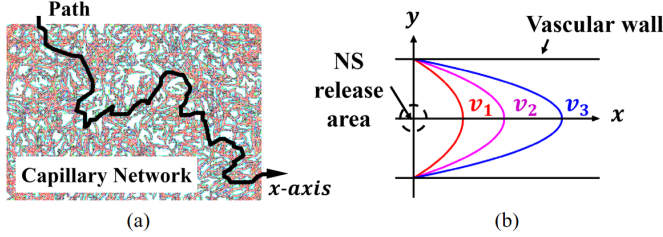


Fig. 2: (a) The NS path inside the capillary network, and (b) the laminar blood flow velocity ($v_3 > v_2 > v_1$) along x -axis.

II. SYSTEM MODELS AND SENSING METHOD

This section presents a framework of *in vivo* sensing of blood viscosity by tracking NS. The flow chart of this framework is shown in Fig. 1. In the forward process, the NS motion is determined by the Brownian noise and the blood flow velocity correlated to the viscosity. In the inverse process, the relative blood viscosity is deduced by monitoring the trajectory of NS with the real-time LQR controller. Furthermore, we estimate the distribution of the location-dependent SNR inside vessel for the control method. The relative displacement of NS is used as feedback for the controller to reduce the noise induced by the Brownian motion, so as to improve the sensing performance.

A. Mathematical Modelling of Blood Flow Velocity

Given that the radius of human vessel is far less than the length of the NS trajectory, we constitute a principal x -axis along the tortuous trajectory of NS in the capillary network as shown in Fig. 2(a). The NS released site is the start point of the x -axis and the y -axis denotes the width of the blood vessel. The blood flow velocity is obtained by solving the Navier-Stokes equation, which provides a fundamental description of laminar flow velocity relating the pressure drop along the vessel and viscosity:

$$V_f(y) = \frac{1}{\eta} \frac{\Delta P}{L} (R^2 - y^2), \quad (1)$$

where $V_f(y)$ is the laminar flow velocity and y is the transversal distance from the x -axis, R is the radius of the

vessel, ΔP is the pressure drop along the vessel with length L , and η denotes the blood viscosity. For a Newtonian fluid, the velocity profile is a parabola with the maximum value located at the center of the vessel, as shown in Fig. 2(b).

B. Kinematic Model of NS

The NS state vector $\mathbf{X}(t)$ is defined as a two-dimensional value $[x, \dot{x}, y, \dot{y}]^T$, where x and \dot{x} are the position and velocity respectively. The following equation illustrates the state-space relation with the mass m :

$$\begin{bmatrix} \dot{x} \\ \ddot{x} \\ \dot{y} \\ \ddot{y} \end{bmatrix} = \mathbf{A} \begin{bmatrix} x \\ \dot{x} \\ y \\ \dot{y} \end{bmatrix} + \mathbf{B} \begin{bmatrix} u_x \\ u_y \end{bmatrix} + \mathbf{W}\varepsilon(t), \quad (2)$$

where

$$\mathbf{A} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{bmatrix}, \quad \mathbf{B} = \begin{bmatrix} 0 & 0 \\ \frac{1}{m} & 0 \\ 0 & 0 \\ 0 & \frac{1}{m} \end{bmatrix}. \quad (3)$$

The global force $[u_x, u_y]^T$ denotes the input, which includes the total affections of the blood flow and the magnetic control. The Brownian noise term $\mathbf{W}\varepsilon(t)$ is a random perturbation produced with magnitude \mathbf{W} . The Brownian force is modelled as a Gaussian random process, thus, it is possible to estimate the NS motion with free diffusion in discrete time. The drag force of blood flow F_d is modelled by the Stokes law:

$$\begin{cases} \gamma = 6\pi\eta a, \\ F_d = \gamma V_f, \end{cases} \quad (4)$$

where the friction coefficient γ is proportional to the blood viscosity η as the NS radius a is setting to a constant value. The combination of the drag force F_d and the Brownian force F_B determines the acceleration and velocity of NS, which is illustrated by the generalized Langevin equations:

$$\tau_B = \frac{m}{\gamma} \approx 10^{-3} s, \quad (5)$$

and

$$\begin{cases} V_x'(t) + \frac{\gamma}{m} V_x(t) = \frac{\gamma V_f}{m} + \frac{F_{B-x}(t)}{m}, \\ V_y'(t) + \frac{\gamma}{m} V_y(t) = \frac{F_{B-y}(t)}{m}, \end{cases} \quad (6)$$

where τ_B is the Brownian timescale for the relaxation of the NS, $V_x(t)$ and $V_y(t)$ are the NS velocity components along the x -axis and y -axis respectively, while $F_{B-x}(t)$ and $F_{B-y}(t)$ are the Brownian force components respectively. In general, the Brownian force is a stochastic impact of the NS performed by the surrounding molecules. The effect of the fluctuating force is given by the first and second moments:

$$\langle F_B(t) \rangle = 0, \quad \langle F_B(t_1), F_B(t_2) \rangle = 2\gamma K_B T \delta(t_1 - t_2), \quad (7)$$

where $\langle \cdot \rangle$ is the average value of stochastic variables, K_B is the Boltzmann's constant and T is the temperature. The delta function in time indicates that there is no correlation

between two time instants t_1 and t_2 . In the simulation, we use a Gaussian random variable following (7) to emulate the Brownian force acting on the NS.

C. Motion-based Viscosity Sensing

Since the kinematic process of NS in the laminar flow follows the generalized Langevin equation, the measurement of blood viscosity can be reduced to a solution of a partial differential equation by removing the Brownian effects. We obtain the following formulas to calculate the blood viscosity, referring to (1), (4) and (6):

$$\begin{cases} \eta = \frac{3\frac{\Delta P}{L}\pi a(R^2 - y^2) + 2F_{B-x} - 2mV_x}{12\pi aV_x}, \\ E(\eta) = \frac{3\frac{\Delta P}{L}\pi a(R^2 - E(y^2)) - 2mE(V'_x)}{12\pi aE(V_x)}, \end{cases} \quad (8)$$

where $E(\cdot)$ is an average operator with respect to the distribution of the variables. As mentioned above, the mean Brownian force should be zero. Furthermore, there is no correlation between Brownian effects in any distinct times. Thus, the Brownian noise can be filtered by analyzing the average trajectory of NS.

In this modelling, the NS are released individually. The blood viscosity is assumed as a time-invariant value, and the parameters apart from the position and velocity in (8) are considered as constant values. Thus, it is possible to obtain the injective relation of different viscosity at different locations. The initial blood viscosity value at the released site is setting as a unit amount. The relative viscosity, calculated by the average trajectory of NS groups at different times, is the ratio of the location-variant viscosity to the initial value.

III. LQR CONTROLLER FOR SENSING ENHANCEMENT

A. Location-dependent SNR

As mentioned above, the blood viscosity is assumed as a time-invariant value. We define blood flow effects as the signal and Brownian motion as the noise, then the location-dependent SNR is represented as

$$SNR = 10\log_{10} \frac{|\gamma V_f(y)V_x|}{|F_B \cdot V_x|}, \quad (9)$$

where $|\gamma V_f(y)V_x|$ is the power of the blood flow and $|F_B \cdot V_x|$ is the power of the Brownian motion along the x -axis. As shown in Fig. 3, the sensing performance declines as the transversal distance increases. Given that a higher SNR along the path is related to a better sensing condition, this motivates us to design the controller stated in the following subsection.

B. Control Method

We propose a real-time LQR-based automatic control method to ensure that the NS trajectory is as close to the centre of the vessel as possible. Without Brownian noise, the controllability matrix is defined as:

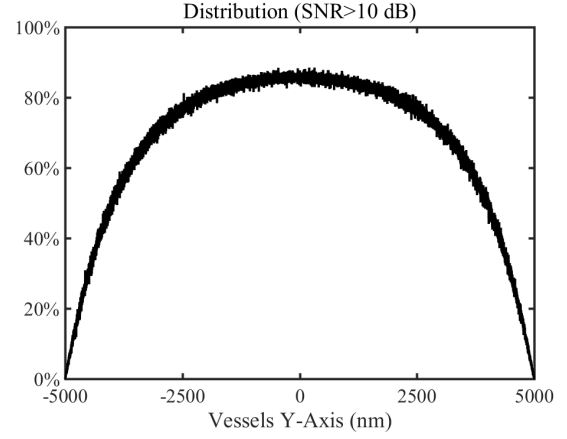


Fig. 3: The distribution of the location-dependent SNR in the vessel.

$$\begin{aligned} \mathbf{C} &= [\mathbf{B}, \mathbf{AB}, \mathbf{A}^2\mathbf{B}, \mathbf{A}^3\mathbf{B}] \\ &= \begin{bmatrix} 0 & 0 & \frac{1}{m} & 0 & 0 & 0 & 0 & 0 \\ \frac{1}{m} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{m} & 0 & 0 & 0 & 0 \\ 0 & \frac{1}{m} & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}, \end{aligned} \quad (10)$$

The full rank of (10) illustrates that a noiseless system is controllable, which is presented by the state vectors of a single NS, $[x, v_x, y, v_y]^T$. Given that the Brownian noise is a time-independent Gaussian White noise, it is hard to design a complete positional controller. However, the magnitude order of magnetic force is much larger than that of Brownian motion. Thus, the impact of the Brownian force is reduced significantly by implementing a high-frequency external magnetic control. Here, LQR is introduced as an optimal controller because of its robustness and cost-effectiveness. Our target is to estimate the optimal control vector $\mathbf{u}^*(t)$, which steers the NS along the expected trajectory. According to [8], the performance function is defined as below:

$$\begin{aligned} J(t) &= \frac{1}{2} \mathbf{x}^T(t_f) \mathbf{S} \mathbf{x}(t_f) + \\ &\frac{1}{2} \int_{t_0}^{t_f} \mathbf{x}^T(t) \mathbf{Q} \mathbf{x}(t) + \mathbf{u}^T(t) \mathbf{R} \mathbf{u}(t) dt, \end{aligned} \quad (11)$$

where t_0 is the initial time and t_f is the end time. And the positive definite coefficient matrices \mathbf{S} , \mathbf{Q} and \mathbf{R} are impacts on the different penalty functions as mentioned below. $L_1 = \frac{1}{2} \mathbf{x}^T(t_f) \mathbf{S} \mathbf{x}(t_f)$ represents the cost function when the controller terminates, $L_2 = \mathbf{x}^T(t) \mathbf{Q} \mathbf{x}(t)$ is the error cost function, and $L_3 = \mathbf{u}^T(t) \mathbf{R} \mathbf{u}(t)$ is the boundary condition of magnetic force.

With parameter-tuning of the above matrixes, the optimal control vector $\mathbf{u}^*(t)$ that minimizes $J(t)$ can be calculated by solving the Riccati equation [8]:

$$\dot{\mathbf{P}}(t) = -\mathbf{P}(t)\mathbf{A} - \mathbf{A}^T\mathbf{P}(t) - \mathbf{Q} + \mathbf{P}(t)\mathbf{B}\mathbf{R}(t)^{-1}\mathbf{B}^T\mathbf{P}(t), \quad (12)$$

and

$$\mathbf{u}^*(t) = -\mathbf{R}(t)^{-1}\mathbf{B}^T\mathbf{P}(t)\mathbf{x}^*(t), t \in [t_0, t_f], \quad (13)$$

where $\mathbf{P}(t)$ is a $n \times n$ positive semi-definite matrix which contributes in solving process.

TABLE I: THE SIMULATION PARAMETERS

Parameters	Value
Blood viscosity (mPa/s)	1.2~2.8
Vessels diameter (μm)	10
Length of sensing path (cm)	2
x -axis unit length (μm)	1
y -axis unit length (nm)	1
Unit discrete time(ms)	0.5
Frequency of controller(Hz)	200
Number of particles	5

IV. SIMULATION RESULTS

In this section, the referential blood viscosity is setting as a benchmark, and the relative viscosity calculated by (8) is compared in two groups. The first group is in the control of external magnetic force while the second group is not. We programmed the computational process on Matlab by choosing some clinical measurement data as the simulation parameters are listed in Tabel I. As shown in Fig. 4. It is obvious that the NS trajectory with a controller is close to the centre of the vessel where the location-dependent SNR is more than 10dB. While without external control, it is a high probability that the NS would adhere to the vessel wall by the accumulated Brownian force. As shown in Fig. 5, the error of the relative viscosity is defined in dB:

$$Err(x) = 10\log_{10} \left| \frac{Vis(x) - Vis^*(x)}{Vis^*(x)} \right|, \quad (14)$$

where $Vis^*(x)$ is the true blood viscosity equals the initial viscosity value in simulation and $Vis(x)$ is estimated relative viscosities.

We can see that the error of the sensing strategy is less than -20dB which means the difference between estimated relative viscosity and benchmark is less than 1%. Due to NS adherence to the vessel wall without the external controller, a sharp and peak value occurs, shown in Fig. 5(a).

V. CONCLUSION

We have proposed a novel sensing strategy to estimate the *in vivo* blood viscosity by tracking the manipulable NS; We have demonstrated through simulating that the proposed strategy can take measurements of the relative blood viscosity accurately under the external controller. Future work may include examining further the impact of non-idealities, such as the blood turbulence, imprecise tracking, and constrained steering. Finally, the proposed sensing strategy should be validated by real experiments to justify further the clinical relevance of the proposed strategy.

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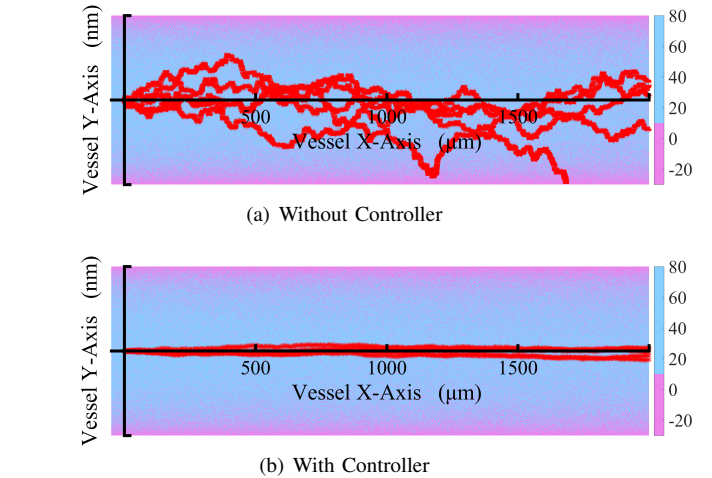


Fig. 4: The trajectories of 5 NS: (a) without controller, and (b) with controller. And the colorbar is the location-dependent SNR (dB).

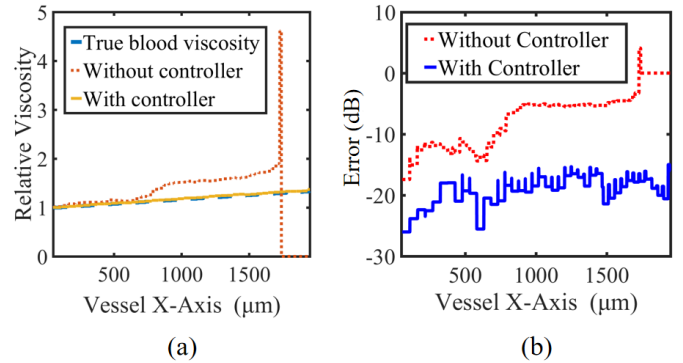


Fig. 5: The sensing results: (a) relative viscosity along the x -axis, and (b) relative error, $Err(x)$.

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